

Adaptive Dynamic Multiscale Methods

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Summary

The aim of this project is the development of a dynamic multiscale approach for the numerical solution of the compressible instationary Euler equations on networks. The modelling is based on the one-dimensional Euler equations. With physically simplified models these equations form a model hierarchy in a natural manner. Using adjoint-based a posteriori error estimators, on accuracy-controlled adaptive simulation of practically relevant gas networks is realized. Core points are an efficient control of spatial and temporal discretizations as well as the model hierarchy to create a basis for multilevel optimization.

Euler Equations

▷ The temperature-dependent instationary Euler equations are

$$\begin{cases} U_t + F(U) &= G(U), \quad (x, t) \in [a, b] \times (0, T], \quad 0 < T, \quad a < b \\ U(x, 0) &= U_0(x), \quad x \in [a, b] \end{cases}$$

with

$$U \equiv \begin{pmatrix} \rho \\ q \\ E \end{pmatrix}, \quad F(U) \equiv \begin{pmatrix} q \\ \frac{q^2}{\rho} + p \\ \frac{q}{\rho}(E + p) \end{pmatrix}, \quad G(U) \equiv \begin{pmatrix} 0 \\ -\frac{\lambda q |q|}{2D\rho} \\ -\frac{k_w}{D}(T - T_w) \end{pmatrix} \quad (1)$$

▷ The variables ρ, q, E, p, T stand for *density, mass flux, total energy, pressure* and *temperature* respectively. It holds $E = \frac{p}{\gamma-1} + \frac{q^2}{\rho}$ and $T = \frac{p}{R\rho}$ to close the system.
 ▷ Further definitions: γ *adiabatic coefficient*, λ *friction coefficient*, D *pipe diameter*, k_w *heat exchange rate*, T_w *outer temperature*, R *gas constant*.

Finite Volume Approach

Integration of (1) over the intervals $I_i \equiv [x_{i-\frac{1}{2}}, x_{i+\frac{1}{2}}]$: $[a, b] = \bigcup_{i=1}^M I_i$ leads to the semi-discrete ODE-system

$$\partial_t \bar{U}_i(t) = \mathcal{L}_i(\mathbf{U})(t) \equiv -\frac{1}{\Delta x} \left(\hat{F}_{i+\frac{1}{2}}(t) - \hat{F}_{i-\frac{1}{2}}(t) \right) + G_i(t), \quad i = 1, \dots, M, \quad (2)$$

with

- ▷ Finite volume $\mathbf{U}(t) \equiv (\bar{U}_i(t))_{i=1}^M$, $\bar{U}_i(t) \equiv \frac{1}{\Delta x} \int_{I_i} U(x, t) dx$,
- ▷ Numerical flux function $\hat{F}_{i+\frac{1}{2}}(\mathbf{U})(t) \approx F(U(x_{i+\frac{1}{2}}))$
- ▷ Numerical integration of the source $G_i(t) \approx \frac{1}{\Delta x} \int_{I_i} G(U(x, t)) dx$
- ▷ Spatial grid size $\Delta x \equiv x_{i+\frac{1}{2}} - x_{i-\frac{1}{2}}$,

Numerical Treatment of the Euler Equations

▷ Involving the *weighted essentially non-oscillatory* procedure \mathcal{W}

$$\partial_t \mathbf{U}(t) = \mathcal{L}(\mathcal{W}(\mathbf{U}))(t) \quad (3)$$

increases the spatial order, cp. [1].

▷ (3) is then solved numerically on the time grid

$$[t^n, t^{n+1}]_{n=1}^N: \bigcup_{n=1}^N [t^n, t^{n+1}], \quad \text{with } \Delta t \equiv t^{n+1} - t^n$$

by an s -stage *singly diagonally implicit Runge-Kutta* method of order 2 (SDIRK(2, s))

$$\begin{cases} Y_j &= \mathbf{U}^n + \Delta t \sum_{k=1}^s a_{j,k} \mathcal{L}(\mathcal{W}(Y_k)), \quad 1 \leq j \leq s \\ \mathbf{U}^{n+1} &= \mathbf{U}^n + \Delta t \sum_{k=1}^s b_k \mathcal{L}(\mathcal{W}(Y_k)) \end{cases} \quad (4)$$

with $\mathbf{U}^n \approx \mathbf{U}(t^n)$ and the optimal parameters

$$a_{k,k} = \frac{1}{2s}, \quad a_{j,k} = \frac{1}{2s}, \quad j < k, \quad b_k = \frac{1}{s}.$$

▷ Whenever $(\mathcal{L}, \|\cdot\|)$ has a contractive *forward Euler* step (FE)

$$\|\mathbf{U} + \Delta_{FE} \mathcal{L}(\mathbf{U})\| \leq \|\mathbf{U}\|$$

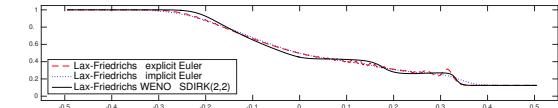
(4) provides a *strong stability preserving* solution in $\|\cdot\|$, cp. [2]

$$\|Y_j\| \leq \|\mathbf{U}^n\|, \quad \forall \Delta t \leq 2s\Delta t_{FE}, \quad 1 \leq j \leq s.$$

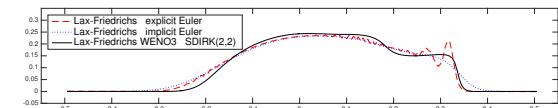
Contribution to Demonstrator 3

Simulation of temperature dependent Euler equations

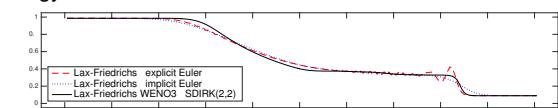
▷ Density



▷ Mass flux



▷ Energy



References

- [1] G. Jiang, C. Shu. Efficient implementation of weighted ENO schemes. Technical report, DTIC Document, 1995
- [2] M.N. Spijker. Stepsize conditions for general monotonicity in numerical initial value problems. *SIAM Journal on Numerical Analysis*, 45(3): 1226-1245, 2007